



Integrated semi groups and some nonlinear fractional integral equations with respect to functions

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Abstract

Some classes of nonlinear fractional abstract integral equations with respect to functions are studied.

Closed operators with α - integrated semi groups are considered.

The existence and uniqueness of solutions of the considered equations are established.

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I. Introduction

Consider the following abstract fractional integral equations with respect to functions:

$$u(t) = u_0 + \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} Au(s) ds + \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} f(s, L(s)u(s)) ds$$

Where $0 < \beta \leq 1$, A is a linear closed operator defined on a dense set S in a Banach space E , and ψ is a real bijective functions, which has continuous derivative $\frac{d\psi(s)}{ds}$ on a finite closed interval $J = [0, T]$, $\psi(0) = 0$,

$\psi(t) \geq 0$, $\frac{d\psi(t)}{dt} > 0$ on J , $\Gamma(\cdot)$

Is the gamma function, $L(t)u(t) = (B_1(t)u(t), \dots, B_r(t)u(t))$, $B_1(t), \dots, B_r(t)$ are families of linear closed operators, defined on dense sets $S_1, \dots, S_r \supset S$ respectively in E . f is a given abstract function defined on $J \times E^r$ to E .

It is assumed that A generates α - times integrated semi groups $\{Q(t) : t \in J\}$, $0 < \alpha \leq 1$ such that $\{Q(t) : t \in J\}$

Is a family of linear bounded operators on E to E , with the following properties :

(i) $Q(t)$ is strongly continuous on J .

(ii) The operator $(\lambda I - A)^{-1}$ exists and:

$$(\lambda I - A)^{-1} = \lambda^\alpha \int_0^\infty e^{-\lambda t} Q(t) dt \quad (1.2)$$

for all $\lambda > \lambda_0 > 0$, the interval (λ_0, ∞) is contained in the resolvent of A , $Q(t)h \in S$, for every $t > 0$, $h \in$

E .

$$(iii) Q(t)h = \frac{t^\alpha}{\Gamma(\alpha+1)} h + \int_0^t A Q(s) h ds \text{ for every } h \in E . \quad (1.3)$$

In section 2, we shall study the linear case.

In section 3, we solve the equation (1.1).

In section 4, we give an example.

The applications can be founded in the theory of elasticity and also in the quantum mechanics, see [1-3].

II. The linear case

Let us study now equation (1.1), when f depends only on t .

In other words, let us try to solve the following equation:

$$u(t) = u_0 + \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - (s)\psi]^{\beta-1} \frac{d\psi(s)}{ds} Au(s) ds + \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - (s)\psi]^{\beta-1} \frac{d\psi(s)}{ds} f(s) ds , \quad (2,1)$$

Suppose that the abstract derivative $\frac{df}{dt}$ Where f is a given abstract continuous function on J , with values in E . exists and continuous on J .

We shall consider the following operators:

$$\Lambda(t) = \int_0^\infty \zeta_\beta(\theta) Q(t^\beta \theta) d\theta$$

$$\Lambda^*(t) = \beta t^{\beta-1} \int_0^\infty \theta \zeta_\beta(\theta) Q(t^\beta \theta) d\theta$$

Where $\zeta_\beta(t)$ is a probability density function defined on $[0, \infty]$ by

$$\zeta_\beta(t) = \frac{1}{\beta} t^{-1-\frac{1}{\beta}} \rho_\beta \left(t^{-\frac{1}{\beta}} \right) \quad (2,2)$$

ρ_β is the one-sided stable probability density function.

The Laplace transform of these functions are given:

$$\int_0^\infty e^{-pt} \rho_\beta(t) dt = \exp(-p^\beta) \quad (2,3)$$

$$\int_0^\infty e^{-pt} \zeta_\beta(t) (t) dt = \sum_{j=0}^\infty \frac{(-p)^j}{\Gamma(1 + \beta j)}$$

We shall consider the following definitions:

$$\begin{aligned} \frac{d^\alpha f(t)}{dt^\alpha} &= \frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_0^t (t-s)^{-\alpha} f(s) ds \\ &= \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-s)^{-\alpha} \frac{df(s)}{ds} ds + \frac{t^{-\alpha}}{\Gamma(1-\alpha)} f(0) \\ I^\alpha f(t) &= \frac{1}{\Gamma(\alpha)} \int_0^t (t-s)^{\alpha-1} f(s) ds \quad 0 < \alpha < 1 \end{aligned}$$

(see [4-12].

Theorem 1.1.

The solution $u(t)$ of equation (2,1) is given by $u(t) = v(\psi(t))$,where

$$v(t) = \frac{a^{\alpha\beta}}{dt^{\alpha\beta}} \left[\Lambda(t)u_0 + \int_0^t \Lambda^*(t - \eta)f(\eta)d\eta \right]$$

Proof.

Consider the equation:

$$v(t) = u_0 + \frac{1}{\Gamma(\beta)} \int_0^t (t - s)^{\beta-1} [Av(s) + f(\psi^{-1}(s))] ds$$

Where ψ^{-1} is the inverse function of ψ .

Set $s = (\psi(\tau))$, so :

$$v(t) = u_0 + \frac{1}{\Gamma(\beta)} \int_0^{\psi^{-1}(t)} (t - \psi(\tau))^{\beta-1} [Av(\psi(\tau)) + f\tau] \frac{d\psi(\tau)}{d\tau} d\tau$$

Thus $u(t) = v(\psi(t))$.

Now let us try to find v .

Let $\tilde{v}(p)$ and $\tilde{g}(p)$ be the Laplace transforms of $v(t)$, and $g(t)$ respectively, where $g(t) = f(\psi^{-1}(t))$.

It is easy to find:

$$\tilde{v}(p) = (p^\beta I - A)^{-1} [p^{\beta-1}u_0 + \tilde{g}(p)].$$

Using (1,2) and (2,4), we get

$$\tilde{v}(p) = \int_0^\infty p^{\alpha\beta} e^{-tp^\beta} Q(t) [p^{\beta-1}u_0 + \tilde{g}(p)] dt \quad (2,5)$$

From (2,3), we can write:

$$\exp(-tp^\beta) = \int_0^\infty e^{-p\theta t^{\frac{1}{\beta}}} \rho_\beta(\theta) d\theta \quad (2,6)$$

Differentiating (2,6) with respect to p , we get

$$\exp(-tp^\beta) = p^{1-\beta} \beta^{-1} \int_0^\infty \theta t^{\frac{1}{\beta}-1} e^{-p\theta t^{\frac{1}{\beta}}} \rho_\beta(\theta) d\theta \quad (2,7)$$

From (2,5), (2,6) and (2,7) one gets:

$$\begin{aligned} \tilde{v}(p) = p^{\alpha\beta} \int_0^\infty e^{-pt} \left[\int_0^\infty \rho_\beta(\theta) Q\left(\frac{t^\beta}{\theta^\beta}\right) u_0 d\theta \right] dt \\ + p^{\alpha\beta} \int_0^\infty e^{-pt} \left[\int_0^\infty \beta \theta^{-\beta} t^{\beta-1} \rho_\beta(\theta) Q\left(\frac{t^\beta}{\theta^\beta}\right) \tilde{g}(p) d\theta \right] dt \quad (2,8) \end{aligned}$$

From (2,2) and (2,8), one gets:

$$\tilde{v}(p) = p^{\alpha\beta} \int_0^{\infty} e^{-pt} \left[\int_0^{\infty} \zeta_{\beta}(\theta) Q(t^{\beta}\theta) u_0 d\theta \right] dp + p^{\alpha\beta} \int_0^{\infty} e^{-pt} \left[\int_0^{\infty} \theta \beta t^{\beta-1} \zeta_{\beta}(\theta) Q(t^{\beta}\theta) \tilde{g}(p) d\theta \right] dp \quad (2,9).$$

According to properties of the Laplace transform of fractional derivatives and noticing that $Q(0)$ is the zero element in E , one gets from (2,9):

$$v(t) = \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} \left[\Lambda(t)u_0 + \int_0^t \Lambda^*(t-\eta)g(\eta)d\eta \right] \quad (2,10)$$

Hence the required result.

Noticing that:

$$\frac{d^{\alpha\beta}}{dt^{\alpha\beta}} t^{\alpha\beta} = \Gamma(\alpha\beta + 1), \int_0^{\infty} \theta^{\alpha} \zeta_{\beta}(\theta) d\theta = \frac{\Gamma(\alpha + 1)}{\Gamma(\alpha\beta + 1)}$$

And using (1,3), (2,10), we get:

$$v(t) = u_0 + \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} \int_0^{\infty} \int_0^{\theta t^{\beta}} \zeta_{\beta}(s) A Q(s) u_0 ds d\theta + \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} \int_0^t \Lambda^*(t-\eta)g(\eta)d\eta \quad (2,11)$$

It is easy to rewrite formula (2,11) in the following form:

$$v(t) = u_0 + \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} \int_0^{\infty} \int_0^{\theta t^{\beta}} \zeta_{\beta} A Q(s) u_0 ds + \int_0^t \Lambda^*(t-\eta) \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} g_0(\eta) d\eta + \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} \int_0^t \Lambda^*(t-\eta)g(\eta)d\eta$$

Where $g_0(s) = g(s) - g(0)$ (2,12)

Also, we can rewrite formula (2,12), in the following form

$$v(t) = u_0 + \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} \int_0^t \Lambda^*(t-\eta)[g(0) + Au_0]d\eta + \int_0^t \Lambda^*(t-\eta) \frac{d^{\alpha\beta}}{dt^{\alpha\beta}} g_0(\eta) d\eta \text{ where } u_0 \in S. \quad (2,13)$$

If $(0) = u(0) = \tilde{0}$, $\tilde{0}$ is the zero in E , we can write

$$v(t) = \int_0^t \Lambda^*(t-\eta) \frac{d^{\alpha\beta}}{d\eta^{\alpha\beta}} g(\eta) d\eta$$

In this case, we get:

$$u(t) = v(\psi(t)) = \int_0^{\psi(t)} \Lambda(\psi(t)-\eta) \frac{d^{\alpha\beta}}{d\eta^{\alpha\beta}} g(\eta) d\eta. \quad (2,14)$$

III. Nonlinear integral equations

Consider the following equation:

$$u(t) = \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} Au(s) ds + \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} f^*(s, L(s)u(s)) ds \quad (3,1)$$

Where $f^*(t, L(t)u(t)) = \frac{1}{\Gamma(\alpha\beta)} \int_0^t (t-s)^{\alpha\beta-1} f(s, L(s)u(s)) ds.$

It is assumed that f is uniformly Hölder continuous in $t \in J$, that is:

$$\|f(t_2, v) - f(t_1, v)\| \leq K(t_2 - t_1)^c, \text{ for all } t_2 > t_1, t_1, t_2 \in J \quad (3,2)$$

Where $v \in E^r$ and K, c are positive constants, $c \leq 1$, $\|\cdot\|$ is the norm in E .

It is assumed also that the Lipchitz condition

$$\|f(t, v^*) - f(t, v)\| \leq K \sum_{i=1}^{i=r} \|v_i^* - v_i\| \quad (3,3)$$

Is satisfied for all $v, v^* \in E$, ($v = (v_1, \dots, v_r)$, $v^* = (v_1^*, \dots, v_r^*)$),

Where K is a positive constant.

About the operators $B_1(t), \dots, B_r(t)$, we assumed that functions $B_1(t)h, \dots, B_r(t)h$ are uniformly Holder continuous in J for $h \in \bigcap_{i=1}^r S_i$.

It is assumed also that

$$\|B(t_2)Q(t_1)h\| \leq \frac{K}{t_1^\gamma} \|h\| \quad (3,4)$$

Where K is positive constant, $0 < \gamma < 1$, $t_2 \in J$, $t_1 \in (0, T]$.

We notice that the solution $u(t)$ of equation (8,1) can be represented by

$$u(t) = v(\psi(t)) \quad (3,5)$$

Where

$$v(t) = \frac{1}{\Gamma(\beta)} \int_0^t (t-s)^{\beta-1} Av(s) ds + \frac{1}{\Gamma(\beta)} \int_0^t (t-s)^{\beta-1} f^*(\psi^{-1}(s), L(\psi^{-1}(s))v(\psi^{-1}(s))) ds \quad (3,6)$$

Set $V(t) = f(t, L(t)u(t))$, we can write, by using formula (2,14), the following representation:

$$v(t) = \int_0^t \Lambda^*(t-\eta) V(\psi^{-1}(\eta)) d\eta$$

Thus

$$u(t) = v(\psi(t)) = \int_0^t \Lambda^*(\psi(t) - \psi(\eta)) \frac{d\psi(\eta)}{d\eta} V(\eta) d\eta \quad (3,7)$$

Let $C_E(J)$ be the set of all abstract continuous functions u on J , with values in E .

We define a distance function $d(u, v)$ by

$$d(u, v) = \max_{t \in J} [e^{-\lambda t} \|u(t) - v(t)\|]$$

Where λ is a positive number. It is clear that $(C_E(J), d(u, v))$ is a complete metric space.

We shall solve equation (3.7).

Theorem 3.1. If $u_1, u_2 \in C_E(J)$ are two solutions of equation (3,7), then $u_1(t) = u_2(t)$, for all $t \in J$.

Proof.

According to the conditions (3,3), (3,4), and gets

$$\|V_1(t) - V_2(t)\| \leq$$

$$\leq K \int_0^t [\psi(t) - \psi(\eta)]^{\delta-1} \|V_1(\eta) - V_2(\eta)\| d\eta$$

Where $\delta = \beta(1 - \varepsilon)$, K is a positive constant, $V_i(t) = f(t, L(t)u_i)$, $i = 1, 2$.

According to the properties of the function ψ and the mean value theorem, we can find a positive constant K such that

$$\|V_1(t) - V_2(t)\| \leq K \int_0^t (t - \eta)^{\delta-1} \|V_1(\eta) - V_2(\eta)\| d\eta$$

It is easy to see that

$$\|V_1(t) - V_2(t)\| \leq K\lambda^{1-\delta} d(V_1, V_2) \int_0^{t-\frac{1}{\lambda}} e^{\lambda s} ds + Kd(V_1, V_2) \int_0^t e^{\lambda s} (t-s)^{\delta-1} ds$$

Thus

$$d(V_1, V_2) \leq K \left(\frac{1}{\lambda}\right)^\delta \left(1 + \frac{1}{\delta}\right) d(V_1, V_2)$$

Choosing λ sufficiently large such that $\varepsilon = K \left(\frac{1}{\lambda}\right)^\delta \left(1 + \frac{1}{\delta}\right) < 1$,

We get $d(V_1, V_2) = 0$. Hence the required result.

Theorem 3.2. Equation (3,7) has a unique solution $u \in C_E(J)$.

Proof. Set

$$V_k(t) = f(t, L(t)u_k(t))$$

Thus

$$d(V_{k+1}, V_k) \leq \nu d(V_k, V_{k-1})$$

By induction, we get

$$d(V_{k+1}, V_k) \leq \nu^k d(V_1, V_2)$$

where V_0 is zero approximation, which can be takes the zero element in E . Thus the sequence

$\{V_k(t)\}$ uniformly converges in the space $C_E(J)$ to a continuous abstract function $V(t)$, which satisfies

$$V(t) = f(t, L(t)u(t))$$

Hence the required result. (see [13-22]).

4-Example

Let $p > 1$, $0 < \alpha < \frac{p-1}{p}$. suppose that $L^p[0,1]$ is the set of all measurable functions f such that $\int_0^1 |f(x)|^p dx$ exists.

Define an operator A by:

$$(Af)(x) = -\frac{d(f(x))}{dx} + \frac{\alpha}{x} f(x)$$

The domain of definition S of A is the set of all absolutely continuous functions f defined on $[0,1]$ with $f(0) = 0$, $\frac{df}{dx} \in L^p[0,1]$.

The considered operator A generates the integrated semi group $Q(t)$, where

$$(Q(t)f)(x) = \int_0^t x^\alpha (x-s)^{-\alpha} f(x-s) H(x-s) ds, \quad x \in [0,1]$$

H is the Heaviside function.

Consider the following equations

$$u(x, t) = \varphi(x) + \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} \left\{ -\frac{\partial u(x, s)}{\partial x} + \frac{\alpha}{x} u(x, s) \right\} ds + \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} f(x, s) ds,$$

$$u(x, t) = \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} \left\{ -\frac{\partial u(x, s)}{\partial x} + \frac{\alpha}{x} u(x, s) \right\} ds$$

$$+ \frac{1}{\Gamma(\beta)} \int_0^t [\psi(t) - \psi(s)]^{\beta-1} \frac{d\psi(s)}{ds} f^*(x, s, u(x, s)) ds$$

$$f^*(x, s, u(x, t)) = \frac{1}{\Gamma(\alpha\beta)} \int_0^t (t-s)^{\beta-1} f(x, s, u(x, s)) ds$$

These equations can be solved as in section 3.

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